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The strain energy density approach applied to bonded joints

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Abstract

Adhesively bonded joints can be treated as bi-materials characterized by a stress singularity promoted by both geometrical and constitutive discontinuities, very close to the point where the interface between two elastic solids intersects a traction-free edge. Different approaches have been proposed to predict their both static and fatigue strength. Most of them are based on the stress distribution at the interface between the adherent and the adhesive. More recently, energy density-based criteria have been developed to predict the static and fatigue strength of mono-material notched components, like welded joints. The most promising one is the strain energy density approach in which it is assumed that failure occurs when the strain energy density, averaged over a control volume of critical radius R_c surrounding the singularity point, will reach a critical value. Advantages are different. Among the others, the strain energy density doesn't depend on the singularity order and doesn't require a fine mesh. This contribution is aimed at applying the strain energy density criterion to bonded joints.

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1. Introduction

The use of adhesives as fixing elements has a thousand-year history, but its first use for industrial applications dates back to the 1940s with the introduction of bonded joints in the aeronautical field. Compared to traditional joints – welding, fastening, bolting – bonded joints guarantee less weight, less cost, thermal and electric insulation, absence of

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distortions, less stress concentration effects, absence of galvanic interaction between dissimilar alloys and finally they are the only choice when one or both the parts to be joined are not metallic materials (say, composite). On the other hand, the preparation of adherents must be very meticulous and the environmental interaction (say, working temperature or contact with solvents) must be carefully taken into account.

The need to predict the mechanical strength of a bonded joint leads to analyzing the stress distribution within the joint itself. This study can be done using an analytical or a numerical approach such as the finite element method (FEM). When the domain of investigation is very close to the point where the interface between two elastic solids intersects a traction-free edge (point A in Fig. 1), a stress singularity is captured by both analytical and numerical models.

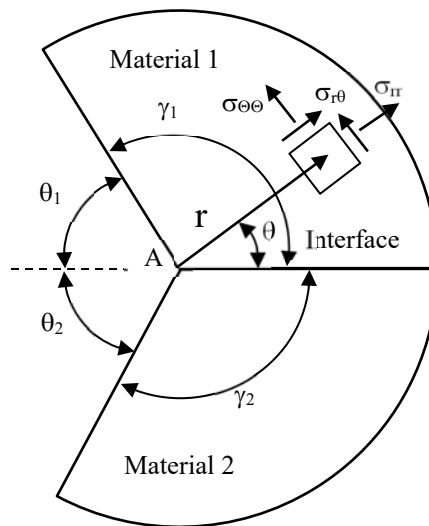


Fig.1. Singularity zone in bi-materials and polar coordinate system

It is worth mentioning that in bi-materials, like bonded joints, the stress singularity is due to both geometrical discontinuities (say, sharp V notches) (Williams, 1952) and constitutive discontinuities at the interface due to the different elastic properties of the two materials coupled.

Over the years, various failure criteria have been developed. The reader can refer to Greenwood et al. (1969), Hart-Smith (1973), Crocombe and Tatarek (1985), Adams and Panes (1994), Ikegami et al. (1989), Crocombe and Adams (1982), Towse et al. (1997a), Towse et al. (1997b), Trantina (1972), Chow and Lu (1992), Zhao, (1991), Schmit and Fraisse (1992). Among these, a local approach based on the intensity of the stress singularity quantified by the stress intensity factor (SIF) has been proposed by Reedy (1993), Reedy and Guess (1993), Lafevbre and Dillard (1999). Unfortunately, SIF parameter requires a very fine mesh near the singularity point and depends on the singularity order, making impossible to compare the strength of bonded joints with different geometries.

In the last years, a powerful approach based on the strain energy density (SED) averaged over a control volume around the singularity has been proposed for static and fatigue design of notched components (Lazzarin, Zambardi, 2001; Berto et al., 2016). It is assumed that the critical radius (R_c) of the control volume is a material characteristic only. The SED criterion was successfully used to quantify the effect of residual stress on high cycle fatigue strength of welded joints, as well (Ferro, 2014; Ferro and Berto, 2016; Ferro et al., 2016). The SED parameter can be calculated with a coarse mesh and due to its scalar nature it is independent from the order of the singularity.

In this work the SED criterion is applied to bonded joints. To this aim, some hypotheses were assumed: the elastic properties of the adherents are much higher than those of the adhesive and the control volume is restricted to the sector belonging to the adhesive only (material 1 in Fig. 1) that will behave as a brittle material. Model parameters are calibrated using experimental data coming from literature for adhesively butt joints made out of steel and epoxy resin.

2. Analytical frame

The strain energy density, under plain stress or plain strain conditions, is given by the following Beltrami's equation:

$$W(r, \theta) = \frac{1}{2E} \left\{ \sigma_r^2 + \sigma_\theta^2 - 2\nu\sigma_r\sigma_\theta + 2(1+\nu)\tau_{r\theta}^2 \right\} \quad (1)$$

where E and ν are the elastic modulus and Poisson's ratio of the adhesive, respectively. On the other hand, the stress distribution near the singularity point is described by (Lazzarin, Zambardi, 2001; Berto et al., 2016):

$$\sigma_{ij}(r, \theta) = \frac{K_1 r^{\lambda-1} f_{ij}(\theta)}{\sqrt{2\pi}} = H_0 r^s f_{ij}^{(0)}(\theta) \quad (2)$$

where K_1 is the mode I SIF, H_0 is the generalized SIF and λ is the eigenvalue (Reedy, 1990). Now, by substituting Eq. (2) in Eq. (1), the following equation holds true:

$$W(r, \theta) = \frac{r^{2(\lambda-1)}}{4\pi E} K_1^2 \left\{ f_{rr}^2(\theta) + f_{\theta\theta}^2(\theta) - 2\nu f_{rr}(\theta) f_{\theta\theta}(\theta) + 2(1+\nu) f_{r\theta}^2(\theta) \right\} = \frac{r^{2(\lambda-1)}}{4\pi E} K_1^2 \beta(\theta) \quad (3)$$

with $\beta(\theta) = f_{rr}^2(\theta) + f_{\theta\theta}^2(\theta) - 2\nu f_{rr}(\theta) f_{\theta\theta}(\theta) + 2(1+\nu) f_{r\theta}^2(\theta)$

Eq. (3) is integrated over the circular sector belonging to the adhesive only, having radius R and $\gamma_1 = \pi/2$ (Fig. 1),

$$E_R = \int_0^R \int_0^{\pi/2} W(r, \theta) dA = \frac{K_1^2}{4\pi E} \int_0^R r^{2(\lambda-1)} r dr \int_0^{\pi/2} \beta(\theta) d\theta = \frac{K_1^2}{4\pi E} \frac{R^{2(\lambda-1)+2}}{2\lambda} \int_0^{\pi/2} \beta(\theta) d\theta \quad (4)$$

and then averaged on the adhesive sector:

$$\bar{W}_R = \frac{E_R}{R^2 \frac{\pi}{4}} = \frac{K_1^2}{\pi^2 E} \frac{R^{2(\lambda-1)}}{2\lambda} \int_0^{\pi/2} \beta(\theta) d\theta = \frac{H_0^2}{\pi E} \frac{R^{2(\lambda-1)}}{\lambda} \int_0^{\pi/2} \beta(\theta) d\theta \quad (5)$$

It is assumed that failure occurs when the strain energy density averaged over a control volume of critical radius R_c surrounding the singularity point will reach a critical value. The R_c value is obtained by equating the energy density to failure of the adhesive with that of the adhesively bonded butt joint:

$$\frac{\sigma_L^2}{2E} = \frac{H_{0,c}^2}{\pi^2 E} \frac{R_c^{2(\lambda-1)}}{\lambda} \int_0^{\pi/2} \beta(\theta) d\theta \quad (6)$$

In Eq. (7), σ_L and $H_{0,c}$ are the stress to failure and the critical GSIF of the adhesive, respectively.

The $H_{0,c}$ value is calculated by combining numerical and experimental results taken from literature (Suzuki, 1985). In particular, the butt joint steel-epoxy (Epikote 828) was taken into account, with the adhesive having the following material properties: $E = 3140$ MPa, $\nu = 0.37$ and $\sigma_L = 65$ MPa. The steel elastic properties are: $E = 206$ GPa, $\nu = 0.3$.

3. Critical radius assessment

The geometry of the adhesively bonded butt joint follows that of Standard ISO 6922. In particular, the joint diameter is 15 mm while the adhesive thicknesses (h) modelled are: 0.12, 0.4, 2.0, 7.5 mm. By taking advantage of the load and geometry symmetry, only a quarter of the geometry has been modelled using PLANE 183 elements and axisymmetric condition. The size of the smallest element was about $1.6 \cdot 10^{-5}$ mm (Fig. 2).

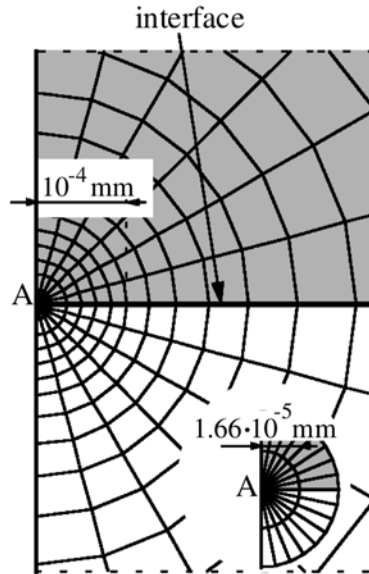


Fig. 2. A detail of the mesh near the singularity point A

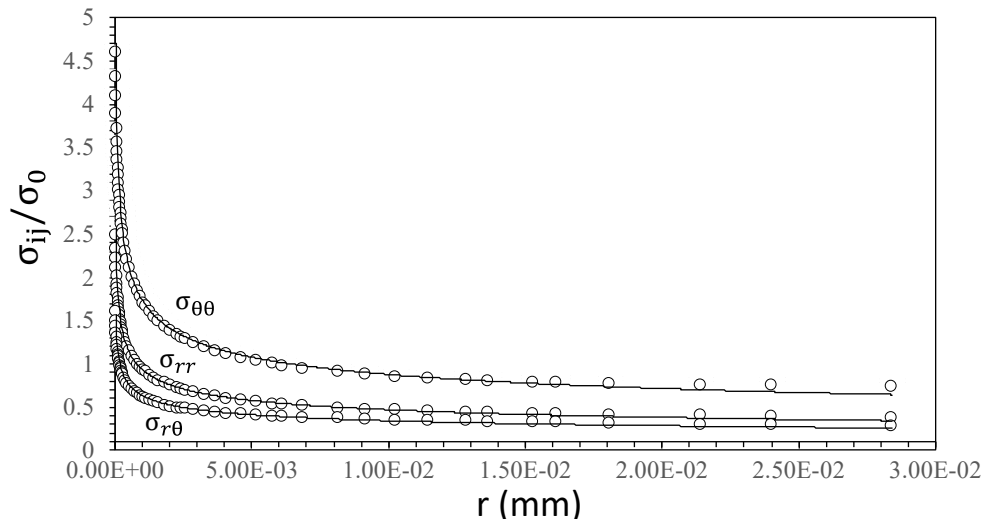


Fig. 3. Asymptotic stress distribution at steel-epoxy interface ($\theta = 0$), adhesive side. $\sigma_0 = 1$ MPa, nominal applied stress; $h = 0.12$ mm

The order of the singularity was found -0.3 ($\lambda = 0.7$), in good agreement with the theoretical value of -0.301 (Akisanya, 1997) (Fig. 3). The GSIF is obtained by using numerical results and Eq. (7):

$$H_0 = \frac{\sigma_{\theta\theta} r^{(\lambda-1)}}{f_{\theta\theta}(0)} = \sigma_{\theta\theta} r^{(\lambda-1)} \quad (7)$$

In particular, by using experimental tensile stress to failure as a function of h , the $H_{0,C}$ value was found to be $10.3 \text{ MPa mm}^{0.3}$, irrespective of adhesive thickness size. It is noted that $f_{\theta\theta}(0)$ is set equal to 1 (Eq. 7) since the angular stress distribution functions are always defined within a constant value. Once these functions are obtained via numerical analysis (Fig. 4), the integral of $\beta(\theta)$ in Eq. (6) is calculated. Finally, by using Eq. (6) the critical radius resulted to be $6.2 \cdot 10^{-3} \text{ mm}$. That low value agrees well with the process zone size of bonded joints (Mintzas and Nowell, 2012; Reedy, 2000), which is much lower than that of welded joints, and is considered an intrinsic property of the bi-material system under investigation, rather than a mere adhesive property only.

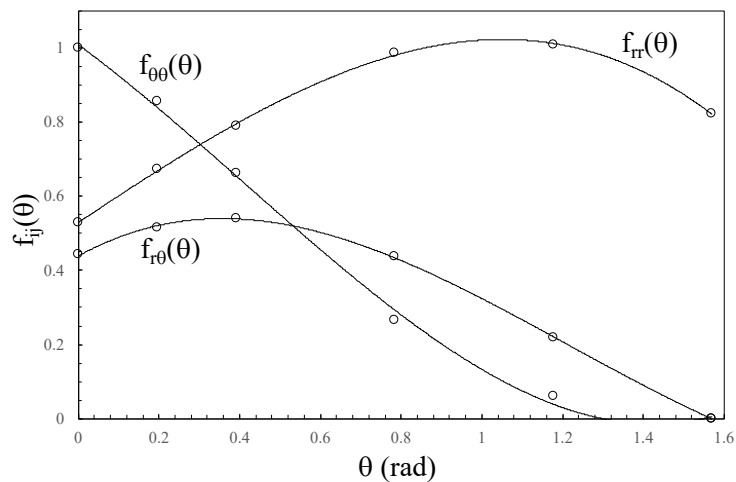


Fig. 4 – Angular stress distribution functions obtained via numerical analysis

4. ASED static failure criterion verification

Table 1 summarizes the results of the simulations in terms of SED values corresponding to the experimental stress to failure as a function of adhesive thickness.

Table 1. Critical SED values derived from experimental results and comparison with the analytical value

Adhesive thickness, h (mm)	Stress to failure σ_1 (MPa)*	$\bar{w}_{r_c}^c$ (mJ/mm ³)**	$\bar{w}_{r_c}^c$ (mJ/mm ³) (Eq. 5)	Δ (%)
0.12	47.1	0.595	0.672	11.4
0.4	31.7	0.592	0.672	11.9
2.0	18.7	0.647	0.672	3.7
7.5	12.1	0.688	0.672	2.3

*(Suzuki, 1985)

**FEM

It is observed that the critical SED value is maintained constant as a function of adhesive thickness, with a maximum deviation from the theoretical value of 11.9% that may be caused by uncertainties on experimental data.

It can be concluded that the SED criterion can be applied to adhesively bonded joint brittle fracture assessment.

5. Conclusions

The strain energy density approach was applied to bonded joints brittle fracture assessment. Under the simplifying assumption of infinitely rigid substrates, the control volume of critical radius R_c is restricted to adhesive only. The R_c value was obtained using experimental data taken from literature concerning the tensile tests carried out on steel/epoxy butt joints. It was found that the SED critical value averaged over the control volume of radius R_c is maintained constant despite the adhesive thickness variation, therefore representing a bi-material property that can be used as failure parameter for adhesively bonded joints design.

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