

# LUNA measurement found no evidence of a low-energy resonance in ${}^6\text{Li}(p, \gamma){}^7\text{Be}$ reaction

D. Piatti<sup>1,2,\*</sup> for LUNA Collaboration

<sup>1</sup>Università degli Studi di Padova, Via Marzolo 8 35136 Padova (Italy)

<sup>2</sup>INFN Division of Padova, Via Marzolo 8 35136 Padova (Italy)

**Abstract.** The  ${}^6\text{Li}(p, \gamma){}^7\text{Be}$  reaction is mainly at work in three nucleosynthesis scenarios: Big Bang Nucleosynthesis,  ${}^6\text{Li}$  depletion in pre-main and in main sequence stars and cosmic ray interaction with interstellar matter. The  ${}^6\text{Li}(p, \gamma){}^7\text{Be}$  S-factor trend was poorly constrained at astrophysical energies because of conflicting experimental results reported in literature. A recent direct measurement, indeed, found a resonance-like structure at  $E_{c.m.} = 195$  keV, corresponding to an excited state at  $E_x \sim 5800$  keV in  ${}^7\text{Be}$  which, however, has not been confirmed by either other direct measurements or predicted by theoretical calculations.

In order to clarify the existence of this resonance, a new experiment was performed at the Laboratory for Underground Nuclear Astrophysics (LUNA), located deep underground in Gran Sasso Laboratory. Thanks to the extremely low background environment, the  ${}^6\text{Li}(p, \gamma){}^7\text{Be}$  cross section was measured in the center-of-mass energy range  $E = 60\text{--}350$  keV with unprecedented sensitivity. No evidence for the alleged resonance was found. LUNA results was confirmed by latest published indirect determination of  ${}^6\text{Li}(p, \gamma){}^7\text{Be}$  S-factor and it is supported by a recent theoretical study.

## 1 Introduction

Lithium abundance involves mainly three nucleosynthesis scenarios. The Galactic chemical evolution models predict, indeed, that most of the solar lithium was provided by low-mass stars [1] while the rest was produced by Big Bang Nucleosynthesis (BBN) [2, 3] or by Galactic cosmic rays interacting with interstellar matter.

The  ${}^6\text{Li}/{}^7\text{Li}$  isotopic ratio has been proposed as a tool to constrain non-standard lithium production mechanisms [4] and pollution of stellar atmospheres [5] in the context of the cosmological lithium problem. Recent (re-)observations of metal poor stars either severely reduced or provided only upper limits for the lithium isotopic ratio [6–8], suggesting that  ${}^6\text{Li}$  depletion must occur in halo stars, which in turn call into question the  ${}^7\text{Li}$  abundance observed in these stars corresponds to the primordial value [9].

The  ${}^6\text{Li}(p, \gamma){}^7\text{Be}$  reaction ( $Q$  value = 5606.85(7) keV) has a crucial role in determining the stellar  ${}^6\text{Li}/{}^7\text{Li}$  ratio. The  ${}^6\text{Li}(p, \gamma){}^7\text{Be}$  reaction not only deplete  ${}^6\text{Li}$  but it also convert some of it to  ${}^7\text{Li}$ , through  ${}^7\text{Be}$  radioactive decay.

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\*e-mail: [denise.piatti@pd.infn.it](mailto:denise.piatti@pd.infn.it)

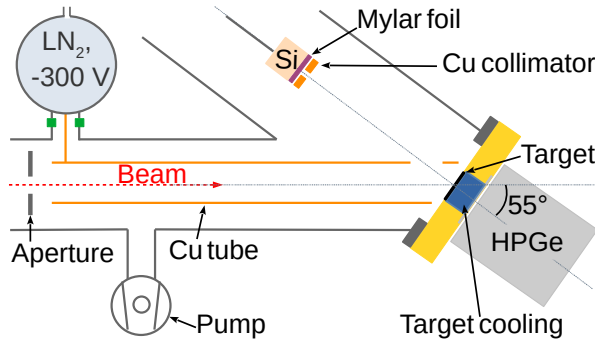
The slope of the astrophysical  $S$ -factor is poorly constrained at low energies given the inconsistent results reported in literature [10, 11]. Moreover, a new resonance at  $E_{c.m.} = 195$  keV, corresponding to an excited level at  $E_x \approx 5800$  keV with  $J^\pi = (1/2^+, 3/2^+)$  and  $\Gamma_p \approx 50$  keV, was claimed by [12]. In a recent comprehensive study of the  ${}^3\text{He}({}^4\text{He},\gamma){}^7\text{Be}$  reaction ( $Q$  value = 1587.14(7)) no evidence for such a resonance was found at  $E_{c.m.} = 4210$  keV [13].

None of the theoretical calculations of the  ${}^6\text{Li}(p,\gamma){}^7\text{Be}$   $S$ -factor can reproduce the newly-reported resonance [14, 15, and references therein], unless this is added *ad-hoc* to reproduce the experimental data [16].

## 2 Experimental Setup

A new experiment [17] was performed at the Laboratory for Underground Nuclear Astrophysics (LUNA), at Laboratori Nazionali del Gran Sasso (Italy) [18]. The LUNA deep-underground location guarantees the reduction of environmental background by several orders of magnitude with respect to overground laboratories, enabling high-sensitivity measurements to be performed.

A schematic view of the experimental setup is shown in Fig.1. The high-intensity pro-



**Figure 1.** Sketch of the experimental setup used for the measurement of the  ${}^6\text{Li}(p,\gamma){}^7\text{Be}$  cross section at LUNA [17].

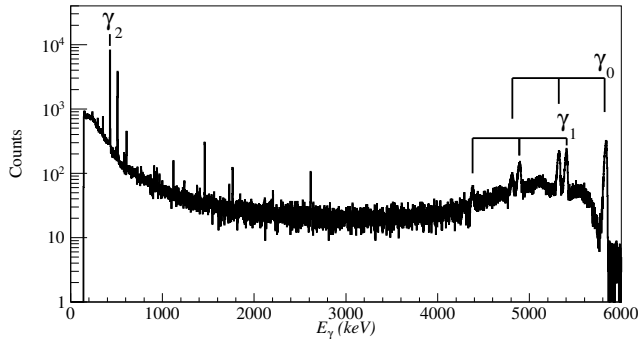
ton beam was provided by LUNA-400 accelerator [19] and it was collimated and delivered through a copper pipe to the target, mounted at  $55^\circ$  with respect to the beam direction. The Cu tube was used both as a cold trap, to improve the scattering chamber vacuum and prevent carbon build-up on target, and for secondary electron suppression. The evaporated targets were made from  ${}^6\text{Li}_2\text{WO}_4$  or  ${}^6\text{Li}_2\text{O}$  powder, with thicknesses  $100 - 200 \mu\text{g}/\text{cm}^2$  and  $20 \mu\text{g}/\text{cm}^2$  respectively. The  ${}^6\text{Li}$  isotopic enrichment level was 95% for all targets, which were water cooled to limit target degradation during irradiation [17].

To detect  ${}^6\text{Li}(p,\gamma){}^7\text{Be}$  reaction  $\gamma$ -rays a High-Purity Germanium (HPGe) detector was positioned in close geometry to the target and at  $55^\circ$  with respect to the beam direction. In addition a Silicon (Si) detector was installed at  $125^\circ$  from the beam direction to detect the  $\alpha$  and  ${}^3\text{He}$  particles from the  ${}^6\text{Li}(p,\alpha){}^3\text{He}$  reaction concurrently with the gamma rays from the  ${}^6\text{Li}(p,\gamma){}^7\text{Be}$  reaction. Efficiencies for both detectors were obtained using GEANT simulations, fine tuned through the comparison with experimental results for  $\gamma$  and  $\alpha$  standards as well as for known resonances of  ${}^{14}\text{N}(p,\gamma){}^{15}\text{O}$  and  ${}^{18}\text{O}(p,\alpha){}^{15}\text{N}$  reactions [17]. The total uncertainty is 4% and 8% for the HPGe and Si detector efficiency respectively.

### 3 Results and Discussion

A measurement of the  ${}^6\text{Li}(p,\gamma){}^7\text{Be}$  and  ${}^6\text{Li}(p,\alpha){}^3\text{He}$  excitation functions was performed for each target in the whole dynamic range of the LUNA-400 accelerator in order to make consistency checks and verify results are unaffected by systematic effects [17].

The  ${}^6\text{Li}(p,\gamma){}^7\text{Be}$  experimental yield was calculated as the sum of the contributions from the direct capture to the ground state ( $\gamma_0$ ) and to the 429 keV excited state of  ${}^7\text{Be}$  ( $\gamma_1$ ).

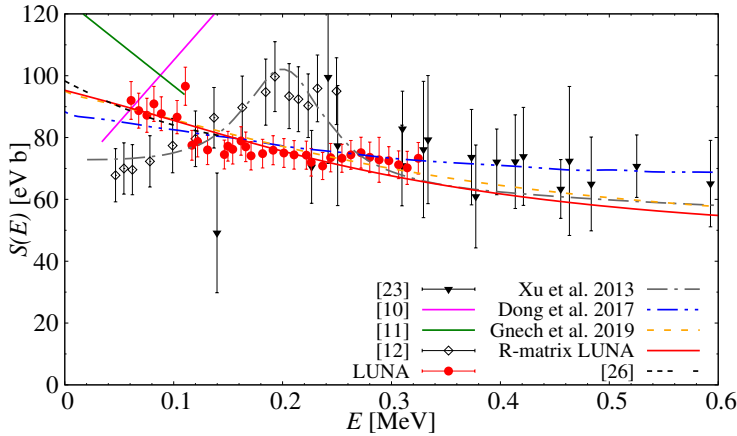


**Figure 2.** Experimental  $\gamma$  spectrum acquired at  $E_p = 265$  keV. The  ${}^6\text{Li}(p,\gamma){}^7\text{Be}$  proceeds through direct capture (DC) to either the ground state of  ${}^7\text{Be}$ ,  $\gamma_0$ , or its first excited state,  $\gamma_1$ , with subsequent emission of a 429 keV secondary gamma ray,  $\gamma_2$ .

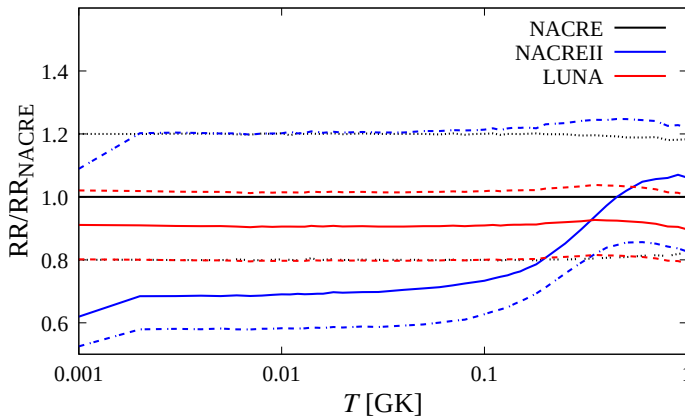
For the calculation of the  ${}^6\text{Li}(p,\gamma){}^7\text{Be}$  reaction  $S$ -factor, we adopted a relative approach [17]: the  $(p,\gamma)$  yield was normalized at each energy to the  $(p,\alpha)$  yield. This ratio can be expressed in terms of the ratio between  $(p,\gamma)$  and  $(p,\alpha)$   $S$ -factors. We adopted for the  ${}^6\text{Li}(p,\alpha){}^3\text{He}$  reaction the  $S$ -factor parametrization reported in [20]. For the  $(p,\alpha)$  channel, the angular distribution coefficients  $A_k$  and related uncertainties were taken from [21, and references therein]. For the  $(p,\gamma)$  channel we adopted the theoretical angular distribution described in [14]. The measured  $S$ -factor was corrected for electron screening using the approximation in [22] and assuming a screening potential  $U_e = 273$  eV [20].

The present  $S$ -factor has a monotonic dependence on the energy and show no evidence of the resonance reported in [12], see Fig.3. The measurement covered the center-of-mass energy range 60 – 350 keV and the reported statistical and systematic uncertainty were  $\leq 2\%$  and 12% respectively. Combining current data and the high energy results reported in [23] an R-matrix fit was performed providing an extrapolated  $S$ -factor to zero energy  $S(0) = 95 \pm 5$  eV b. The R-matrix fit was used to calculate a new  ${}^6\text{Li}(p,\gamma){}^7\text{Be}$  reaction rate, which is 9% lower than NACRE [24] and 33% higher than reported in NACREII [16] at temperatures relevant for  ${}^6\text{Li}$  depletion in pre-main sequence stars. Moreover the reaction rate uncertainty has been significantly reduced [17], see Fig.4.

The result of a subsequent indirect study confirms LUNA extrapolation down to low energies for the  ${}^6\text{Li}(p,\gamma){}^7\text{Be}$   $S$ -factor, reporting an  $S(0) = 92 \pm 12$  eV b [25]. A recent theoretical study found a consistent trend for the  ${}^6\text{Li}(p,\gamma){}^7\text{Be}$   $S$ -factor predicting a  $S$ -factor to zero energy of 98.3 eV b [26]



**Figure 3.** Astrophysical S-factor for the  ${}^6\text{Li}(p,\gamma){}^7\text{Be}$  reaction as obtained by LUNA in red [17]. Previous experimental data and theoretical evaluations are also shown for comparison. The solid red line represents an R-matrix fit of LUNA data and data from [23].



**Figure 4.** Reaction rate for the  ${}^6\text{Li}(p,\gamma){}^7\text{Be}$  reaction, normalized to the NACRE rate [24]. The NACRE II rate [16] is also shown for comparison. Dashed lines represent the uncertainty on the NACRE rate (black), on NACREII rate (blue) and on LUNA rate (red).

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